

THEORY OF THE KENNICUTT-SCHMIDT LAW FOR STAR FORMATION IN GALAXIES

MARK KRUMHOLZ & CHRIS MCKEE

ApJ 630, 250 (2005)

Workshop on the Kennicutt-Schmidt Law

UCSD 12/18-19/2006

REQUIREMENTS FOR A THEORY OF THE KS LAW

Two forms consistent with observation:

$$d\Sigma_*/dt \propto \Sigma_g^{1.4} \quad d\Sigma_*/dt \propto \Sigma_g/t_{\text{dyn}}$$

Applies over 8 orders of magnitude in star formation rate $d\Sigma_*/dt$
from normal galaxies to starbursts

Star formation inefficient: gas depletion time \gg dynamical time t_{dyn}

Previous theories

Padoan (1995): Predicts SFR in turbulent GMCs, but no prescription for application to galaxies in which GMC properties are not observed

Silk (1997): SFR set by supernova feedback; depends on uncertain porosity of hot gas

Tan (2000): SFR set by cloud-cloud collisions; normalization set by comparison with obs.

Kravtsov (2003) & Li et al (2005): simulations show fraction of high-density gas $\propto \Sigma_g^{1.4}$
but definition of “high-density” arbitrary and rate of SFR not determined

Elmegreen (2002, 2003): $\text{SFR}/\text{volume} = \epsilon_{\text{core}} f_c (G\rho_c)^{1/2} \rho$ with $\epsilon_{\text{core}} = 1/2$

Fraction of gas in dense cores, f_c , determined from observed KS law

Corresponding density is $\rho_c/\rho = 10^5$, but it is not clear why this is the critical density, nor how it should vary in different galaxies

TURBULENCE-REGULATED STAR FORMATION (KM05)

Assume (1) that star formation occurs in GMCs

$$\begin{aligned} d\Sigma_*/dt &= \text{SFR}_{\text{ff}} \Sigma_{\text{GMC}} / t_{\text{ff}} \\ &= \text{SFR}_{\text{ff}} f_{\text{GMC}} \Sigma_{\text{g}} / t_{\text{ff}} \end{aligned}$$

where t_{ff} is the free-fall time in a GMC

SFR_{ff} is the fraction of gas that goes into stars per free-fall time

f_{GMC} is the fraction of gas in GMCs

$$\begin{aligned} &\approx (1 + 0.025 / \Sigma_{\text{g},2}^2)^{-1} \text{ from Rosolowsky \& Blitz (2006),} \\ &\text{where } \Sigma_{\text{g},2} = \Sigma_{\text{g}} / 100 M_{\text{sun}} \text{ pc}^{-2} \end{aligned}$$

Objective: Determine SFR_{ff} and t_{ff} in terms of Σ_{g} and t_{dyn}

The Star Formation Rate per Free-Fall Time SFR_{ff}

Assume (2) that the probability distribution $p(x)$ of the density in GMCs is log normal as is appropriate for supersonically turbulent gas (e.g., Padoan & Nordlund 2002)

Let $x = \rho/\rho_0$ where ρ_0 is the average density in the GMC

Then $dp(x) \propto \exp [- (\ln x - \langle \ln x \rangle)^2 / 2 \sigma_\rho^2]$

where $\sigma_\rho^2 \approx \ln (1 + 0.75 \mathcal{M}^2)$ and $\mathcal{M} = \sigma/c_s$ is the Mach number

Assume (3) that gas above some critical density ρ_{cr} forms stars with an efficiency ϵ_{core} at a rate corresponding to some number (ϕ_t) of free-fall times:

$$dM_*/dt = (M \epsilon_{\text{core}} / \phi_t t_{\text{ff}}) \int_{x_{\text{cr}}} x dp(x)$$

where $x_{\text{cr}} = \rho/\rho_{\text{cr}}$

$$\Rightarrow \text{SFR}_{\text{ff}} = (\epsilon_{\text{core}} / \phi_t) \int_{x_{\text{cr}}} x dp(x) \quad (\text{can be evaluated analytically})$$

For numerical evaluation, we take $\epsilon_{\text{core}} \approx 1/2$ (Matzner & McKee 2000)

What is the Minimum Density for Star Formation, ρ_{cr} ?

Assume (3') that stars of average mass form in cores dominated by thermal, not turbulent pressure --- valid in Galactic GMCs

Sonic length: Scale at which turbulent motions match thermal motions:

$$\text{Line-width size relation } \sigma_l = c_s (l / \lambda_s)^{1/2} \Rightarrow$$

$$\lambda_s = 2R (c_s / \sigma_{2R})^2 = 2R / \mathcal{M}^2 \quad \text{where } 2R \text{ is the diameter of the GMC}$$

Gravitational collapse possible in a thermal core only if the sonic length $>$ Jeans length (see Padoan 95; equivalent to mass inside sonic length $>$ \sim Bonnor-Ebert mass)

Can show this implies $\rho_{\text{cr}} = 0.8 \alpha_{\text{vir}} \mathcal{M}^2 \rho_0$

where $\alpha_{\text{vir}} = 5\sigma^2 R / GM \sim 1$ is the virial parameter

Notes: (1) this corresponds to $P_{\text{cr}} = \rho_{\text{cr}} c_s^2 \approx \rho_0 \sigma^2$: the critical thermal pressure is comparable to the turbulent pressure in the GMC (Padoan 95)

(2) low-mass cores have $\alpha_{\text{vir}} \sim \mathcal{M} \sim 1$ and are therefore at the critical density

Evaluation of the Star Formation Rate per Free-Fall Time, SFR_{ff}

Recall $\text{SFR}_{\text{ff}} = (\epsilon_{\text{core}} / \phi_t) \int_{x_{\text{cr}}} x \, dp(x)$

We now know x_{cr} and we adopt $\epsilon_{\text{core}} = 1/2$

Vazquez-Semadeni et al. (2003) carried out hydrodynamic simulations and showed that the star formation rate depends on the sonic length. Fitting to their results gives $\phi_t = 1.9$ as the number of free-fall times (evaluated at ρ_0) required for core collapse.

A power-law fit to our results yields

$$\text{SFR}_{\text{ff}} \approx 0.017 \alpha_{\text{vir}}^{-0.7} (\mathcal{M} / 100)^{-0.3}$$

⇒

Star formation is inefficient (a few percent), in agreement with observation.

The rate depends only weakly on the Mach number \mathcal{M} ; note that $\alpha_{\text{vir}} \sim 1$ in GMCs

New Form of the KS Law

We now have

$$\begin{aligned} d\Sigma_*/dt &= \text{SFR}_{\text{ff}} f_{\text{GMC}} \Sigma_g / t_{\text{ff}} \\ &= 0.017 \alpha_{\text{vir}}^{-0.7} (\mathcal{M}/100)^{-0.3} f_{\text{GMC}} \Sigma_g / t_{\text{ff}} \end{aligned}$$

What are the Mach number \mathcal{M} and the free-fall time $t_{\text{ff}} \propto \rho_0^{-1/2}$?

Assume star-forming disk is marginally stable so that $Q \approx 1$

$$\Rightarrow \sigma_g = \pi G \Sigma_g Q / 2^{1/2} \Omega \quad \text{in disk, where } \Omega \text{ is the angular velocity}$$

Density in the disk is given by $\rho_g = P_g / \sigma_g^2$ where $P_g \approx (\pi/2)G \Sigma_g \Sigma_{\text{tot}}$

Pressure in GMC $\approx (2-10) P_g$ and density in GMC $\approx (2-7) \rho_g \Rightarrow \sigma_{\text{GMC}} \approx \sigma_g$
and $t_{\text{ff}} \propto \rho_0^{-1/2} \propto \Omega$ (details in KM05)

$$\begin{aligned} \Rightarrow d\Sigma_*/dt &\approx 0.16 \mathcal{M}^{-0.3} f_{\text{GMC}} \Sigma_g \Omega \quad \text{similar to KS law except for } \mathcal{M}^{-0.3} \\ &\approx 9.5 f_{\text{GMC}} \Sigma_{g,2}^{0.7} \Omega_6^{1.3} M_{\text{sun}} \text{ yr}^{-1} \text{ kpc}^{-2}, \quad \text{where } \Omega_6 = \Omega \times 10^6 \text{ yr} \end{aligned}$$

Star Formation Threshold

Star formation cuts off in outer regions of galaxies

Generally attributed to Toomre Q rising above 1 \Rightarrow stable

Predicted SFR varies as $Q^{-1.3} f_{\text{GMC}}$: declines rapidly in outer regions

since Q increases and the molecular fraction decreases;

in addition, a smaller fraction of the molecular gas is in GMCs at large radii.

Test of New Form: $d\Sigma_*/dt \approx 9.5 f_{\text{GMC}} \Sigma_{\text{g},2}^{0.7} \Omega_6^{1.3} M_{\text{sun}} \text{yr}^{-1} \text{kpc}^{-2}$

Should apply to individual galaxies as well as sample of galaxies

Does not apply to individual GMCs since expect large fluctuations in the star formation rate (Krumholz, Matzner & McKee 2006)

Milky Way: use f_{GMC} and $\Sigma_{\text{g},2}$ from observation
calculate $Q(r)$ in spiral arms (suppressed in above expression)

Predict SFR between 3 and 11 kpc of $4.5 M_{\text{sun}} \text{yr}^{-1}$

Consistent with observed rate $\approx 3 M_{\text{sun}} \text{yr}^{-1}$ (McKee & Williams 1997)

Comparison with Classical Forms of KS Law

There are two forms of the KS law because Σ_g and Ω are correlated in the data:

$$\Omega_6 \approx 0.06 \Sigma_{g,2}^{0.5} \quad \text{for } \Sigma_g > 1 \text{ M}_{\text{sun}} \text{pc}^{-2}$$

First form:

$$\text{Observed: } d\Sigma_*/dt = 0.16 \Sigma_{g,2}^{1.4} \quad \text{M}_{\text{sun}} \text{yr}^{-1} \text{kpc}^{-2} \quad (\text{Kennicutt 1998})$$

$$\text{Theory: } d\Sigma_*/dt = 0.19 f_{\text{GMC}} \Sigma_{g,2}^{1.3} \quad \text{M}_{\text{sun}} \text{yr}^{-1} \text{kpc}^{-2}$$

Second form:

$$\text{Observed: } d\Sigma_*/dt = 1.7 \Omega_6 \Sigma_{g,2} \quad \text{M}_{\text{sun}} \text{yr}^{-1} \text{kpc}^{-2} \quad (\text{Kennicutt 1998})$$

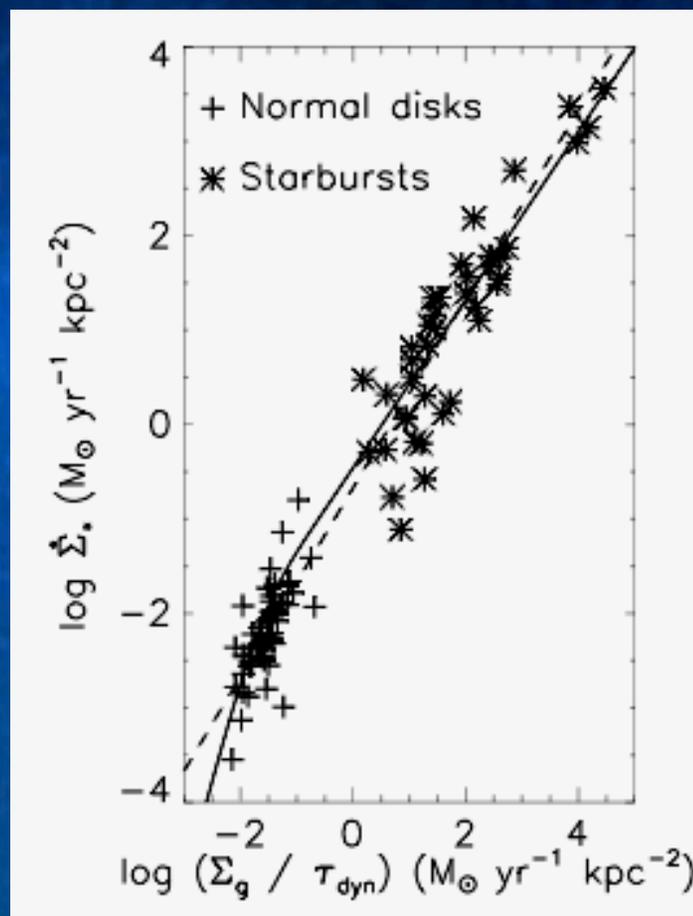
$$\text{Theory: } d\Sigma_*/dt = 3.2 f_{\text{GMC}} (\Omega_6 \Sigma_{g,2})^{0.9} \quad \text{M}_{\text{sun}} \text{yr}^{-1} \text{kpc}^{-2}$$

Comparison with Second Form of KS Law

$$\tau_{\text{dyn}} = 4\pi / \Omega$$

- - - Power-law fit (Kennicutt 98)

— Theory



Tests of Theory

OBSERVATIONAL:

- * Test SFR_{ff} from observations of a sample of GMCs
- * Test $d\Sigma_*/dt$ in annular rings in galaxies
- * Increase the sample size to break the degeneracy between the two forms of the KS law and between observation and theory

THEORETICAL:

- + Predicts time scale for star cluster formation of 3-4 dynamical times, consistent with observation (Tan, Krumholz, & McKee 2006)
- + When used in a dynamical model for GMC evolution, successfully predicts GMC lifetimes and column densities (Krumholz, Matzner, & McKee 2006)

Extending the Theory

- * Determine the GMC fraction f_{GMC} theoretically
Particularly important for low-metallicity galaxies and high-redshift galaxies
- * Determine the effects of magnetic fields
Could alter density PDF and slow rate of star formation
Observations of fields in the Galaxy suggest effects are modest
- * Predict the level of turbulence in GMCs (i.e., predict α_{vir})
Understand the driving mechanisms that counter turbulent decay
Particularly puzzling in starbursts, where σ larger than given by
plausible momentum sources other than self-gravity
- * Show how the massive stars that are observed are related to the low-mass
stars predicted by the theory (i.e., understand the IMF)

These are some of the fundamental questions of star formation

CONCLUSION

The assumptions that

- Stars form in virialized GMCs that are supersonically turbulent
- The density distribution is log normal, as expected for such turbulence in isothermal gas
- Gas dense enough that thermally supported cores that are gravitationally unstable form stars with an efficiency $\epsilon_{\text{core}} \sim 1/2$

imply a star formation law that should apply when averaged over a large number of GMCs, whether in a single galaxy or many galaxies:

$$d\Sigma_{*}/dt \approx 9.5 f_{\text{GMC}} \Sigma_{\text{g},2}^{0.7} \Omega_6^{1.3} M_{\text{sun}} \text{ yr}^{-1} \text{ kpc}^{-2}$$

This result is consistent with existing observations